

Mesons as Open Strings

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with **T. Imoto and T. Sakai**

(in preparation)

1 Introduction

mesons ($N_f = 2$, Isovector)

parity JPC charge conjugation
spin JPC mass (MeV)
 Δ ... not established

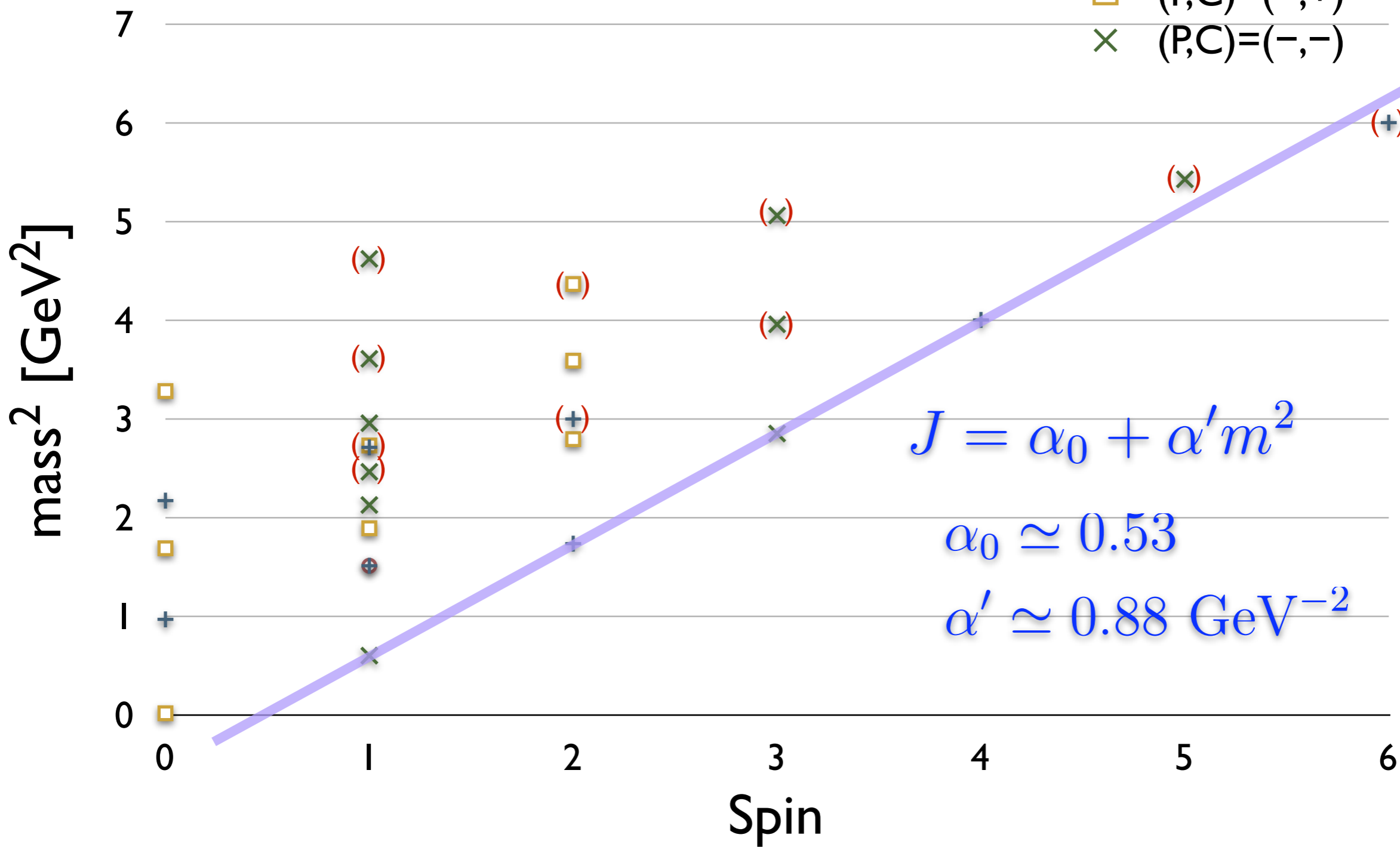
$0^{-+}(\pi)$	135	1300	1812			
$0^{++}(a_0)$	985	1474				
$1^{--}(\rho)$	776	1459	1570^Δ	1720	1900^Δ	2150^Δ
$1^{++}(a_1)$	1230	1647^Δ				
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Q: How can we understand these numbers?



Hint: Regge trajectory

- + (P,C)=(+,+)
- (P,C)=(+,-)
- (P,C)=(-,+)
- × (P,C)=(-,-)



Mesons are strings !?

● **Difficulties in the old days**

- Consistent in 10 dim space-time
- \exists massless particles with $J = 1$ and $J = 2$
(open) (closed)

$$J = \alpha_0 + \alpha' m^2 \quad \alpha_0 = 1 ! \quad (\text{for open string})$$

Not consistent with meson spectrum !?

● **Gauge / String duality suggests**

- 4 dim QCD \longleftrightarrow 10 dim string theory
dual (in a certain curved background)
- massive mesons \longleftrightarrow massless gauge field
in 4 dim dual in 10 dim

The above difficulties can be solved !!

Holographic QCD

- 4 dim QCD \longleftrightarrow 10 dim string theory
(in a certain curved background)
dual

holographic QCD

D4/D8-branes in type IIA string theory

(reviewed in the next section)

[T.Sakai and S.S. 04]

- mesons are open strings on D8
- π, ρ, a_1 , etc. are obtained from the massless mode

\uparrow 0^{-+} \uparrow 1^{--} \uparrow 1^{++}

$$m_{a_1}/m_\rho \simeq \begin{cases} 1.53 & (\text{theory}) \\ 1.59 & (\text{exp}) \end{cases}$$

Q: What about the other mesons ?



Consider massive modes (excited strings)

Plan

- ✓ 1 Introduction
- 2 Brief review of the model
- 3 Meson spectrum
- 4 Comparison with data
- 5 Discussion

2

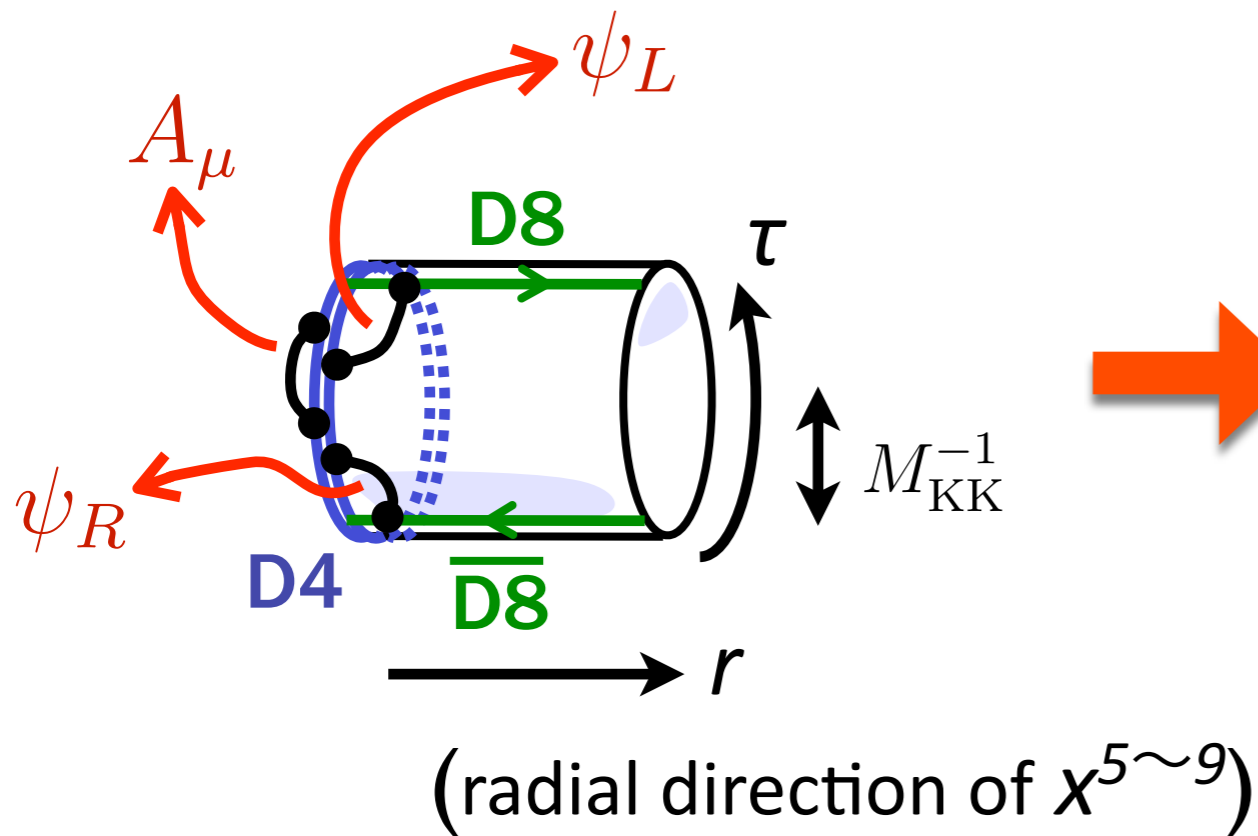
Brief review of the model

[T. Sakai and S.S. 04]

● Brane configuration

	x^0	x^1	x^2	x^3	τ	x^5	x^6	x^7	x^8	x^9
D4 \times N_c	○	○	○	○	○					
D8- $\overline{\text{D8}}$ \times N_f	○	○	○	○		○	○	○	○	○

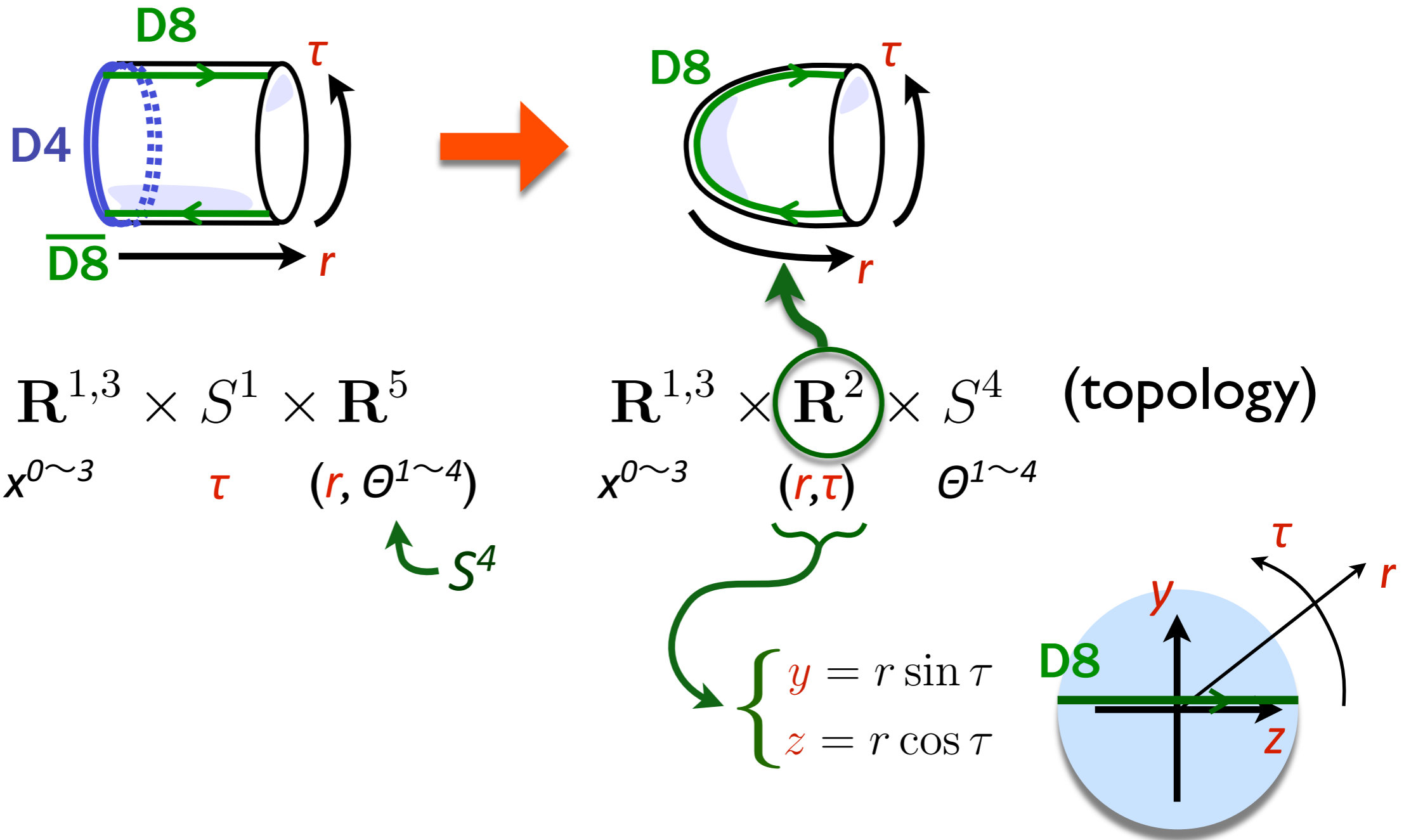
S^1 with ~~SUSY~~ b.c.



4 dim $U(N_c)$ QCD
with N_f massless quarks
(at low energy)

Holographic description

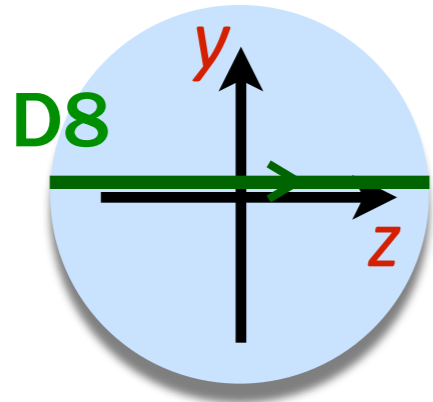
- replace D4 with the corresponding curved background [Witten 98]
- D8 are treated as probe brane (assuming $N_c \gg N_f$)



● Hadrons in the model

$$\mathbf{R}^{1,3} \times \mathbf{R}^2 \times S^4$$

$x^{0\sim 3}$ (z, y)



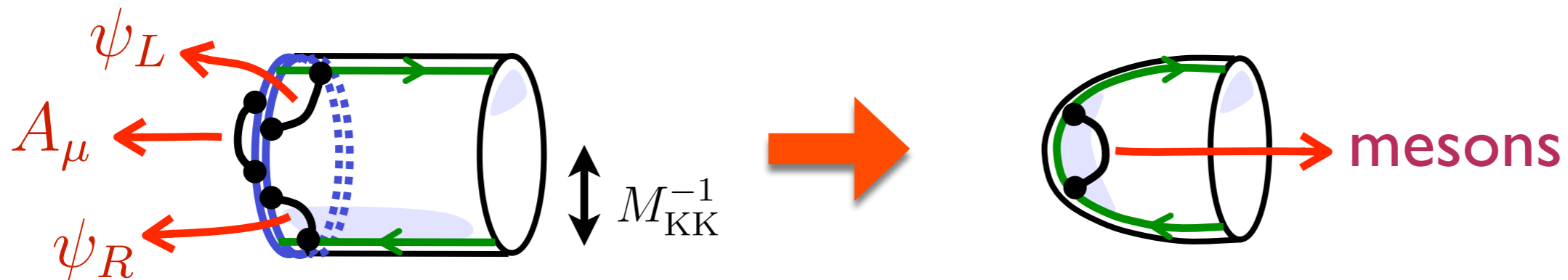
D8 are extended along $\overbrace{(x^\mu, z)}^{5\text{dim}} \times S^4$
 $(\mu=0\sim 3)$

- closed strings \longleftrightarrow glueballs
- open strings on D8 \longleftrightarrow mesons
- D4 wrapped on S^4 \longleftrightarrow baryons

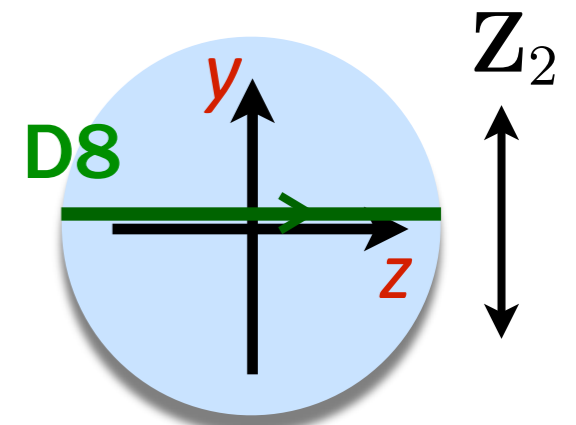
today's
topic

QCD mesons vs artifacts

- Our brane config. is invariant under $SO(5) \curvearrowright S^4$
- quarks and gluons are invariant under $SO(5)$
(non-invariant states are massive KK-modes)



- QCD mesons are $SO(5)$ invariant states
(non-invariant states are KK-artifacts)
- Similarly, we can show that QCD mesons are invariant under Z_2 sym generated by $I_{y9}(-1)^{F_L}$
 $I_{y9}: (y, x^9) \rightarrow (-y, -x^9) \quad (\tau \rightarrow -\tau)$


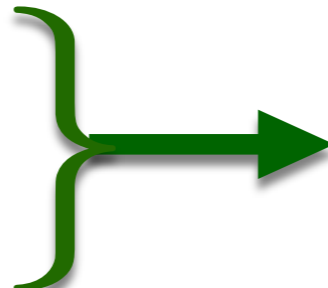



Consider $SO(5) \times Z_2$ invariant states

3 Meson spectrum

Consider open strings attached on D8

● Massless mode (light-cone gauge: $x^\pm = x^0 \pm x^1$)

- $\psi^I_{-1/2}|0\rangle_{\text{NS}}$ ($I = 2, 3, z$)  5 dim gauge field A_μ, A_z
 - $\psi^A_{-1/2}|0\rangle_{\text{NS}}$ ($A = y, 6, 7, 8, 9$) 
 - $|\alpha\rangle_{\text{R}}$ spinor 
- S^4
- not invariant under $SO(5) \times \mathbf{Z}_2$

● KK decomposition along z direction

$$A_\mu(x^\mu, z) = \sum_{n=1}^{\infty} B_\mu^{(n)}(x^\mu) \psi_n(z)$$

$$A_z(x^\mu, z) = \sum_{n=0}^{\infty} \varphi^{(n)}(x^\mu) \phi_n(z)$$

complet sets

	$B_\mu^{(1)}$	$B_\mu^{(2)}$	$B_\mu^{(3)}$...	$\varphi^{(0)}$	$\varphi^{(1)}$...
J^{PC}	1^{--}	1^{++}	1^{--}	...	0^{-+}	eaten	
	ρ	a_1	ρ'	...	π		

[T.Sakai and S.S. 04]

● First excited massive modes

First, we consider flat space-time limit
(justified when 't Hooft coupling is very large)

$SO(5) \times \mathbf{Z}_2$ invariant states:

- $\psi_{-3/2}^I |0\rangle_{\text{NS}}$
 - $\alpha_{-1}^{(I} \psi_{-1/2}^{J)} |0\rangle_{\text{NS}}$
 - $\alpha_{-1}^{[I} \psi_{-1/2}^{J]} |0\rangle_{\text{NS}}$
 - $\psi_{-1/2}^I \psi_{-1/2}^J \psi_{-1/2}^K |0\rangle_{\text{NS}}$
 - $\alpha_{-1}^y \psi_{-1/2}^y |0\rangle_{\text{NS}}$
 - $\sum_{a=6,7,8,9} \alpha_{-1}^a \psi_{-1/2}^a |0\rangle_{\text{NS}}$
- $\underbrace{\hspace{10em}}_{S^4}$

$(I = 2, 3, z)$

5dim field	$SO(4)$ little gr
h_{MN} ($M, N = 1, 2, 3, z$)	
A_{MNP}	
$\varphi^{[1]}$	1
$\varphi^{[2]}$	1

● KK decomposition along z direction

$$h_{MN}(x^\mu, z) = \sum_{n=0}^{\infty} h_{MN}^{(n)}(x^\mu) \phi_n(z) \quad \text{etc.}$$

lowest modes: $(i, j, k = 1, 2, 3)$

	$h_{ij}^{(0)}$	$h_{iz}^{(0)}$	$h_{zz}^{(0)}$	$A_{ijk}^{(0)}$	$A_{ijz}^{(0)}$	$\varphi^{[1,2](0)}$
J^{PC}	2^{++}	1^{+-}	0^{++}	0^{-+}	1^{--}	$0^{++} \times 2$

● Second excited massive mode

lowest modes:

J^{PC}	3^{--}	2^{++}	2^{--}	$2^{-+} \times 2$	$1^{--} \times 7$	$1^{++} \times 3$	$1^{+-} \times 4$	1^{-+}	$0^{++} \times 2$	$0^{-+} \times 6$
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- **Mass formula** (naive shortcut)

- Flat space-time limit:

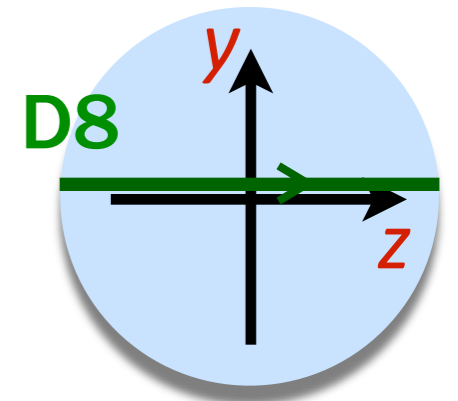
$$m_0^2 = \frac{N}{\alpha'}$$

$$\left(\alpha'^{-1} = \frac{4}{27} \lambda M_{\text{KK}}^2 \right)$$

$N = 0, 1, 2, \dots$: excitation level

- Massive particle in curved space-time

$$S = -m_0 \int \sqrt{g_{tt}} dt = -m_0 (1 + z^2)^{1/4} \int dt$$



→ particle in potential: $V(z) = m_0 (1 + z^2)^{1/4}$

→ $M_n \simeq m_0 + \frac{1}{\sqrt{2}} \left(n + \frac{1}{2} \right) M_{\text{KK}} + \mathcal{O}(\lambda^{-1/2})$ $n = 0, 1, 2, \dots$

harmonic oscillator approx.

- More careful analysis shows that the $O(1)$ term is not affected by the RR-flux, α' correction, etc.

4

Comparison with data

Massless mode

[T.Sakai and S.S. 04]

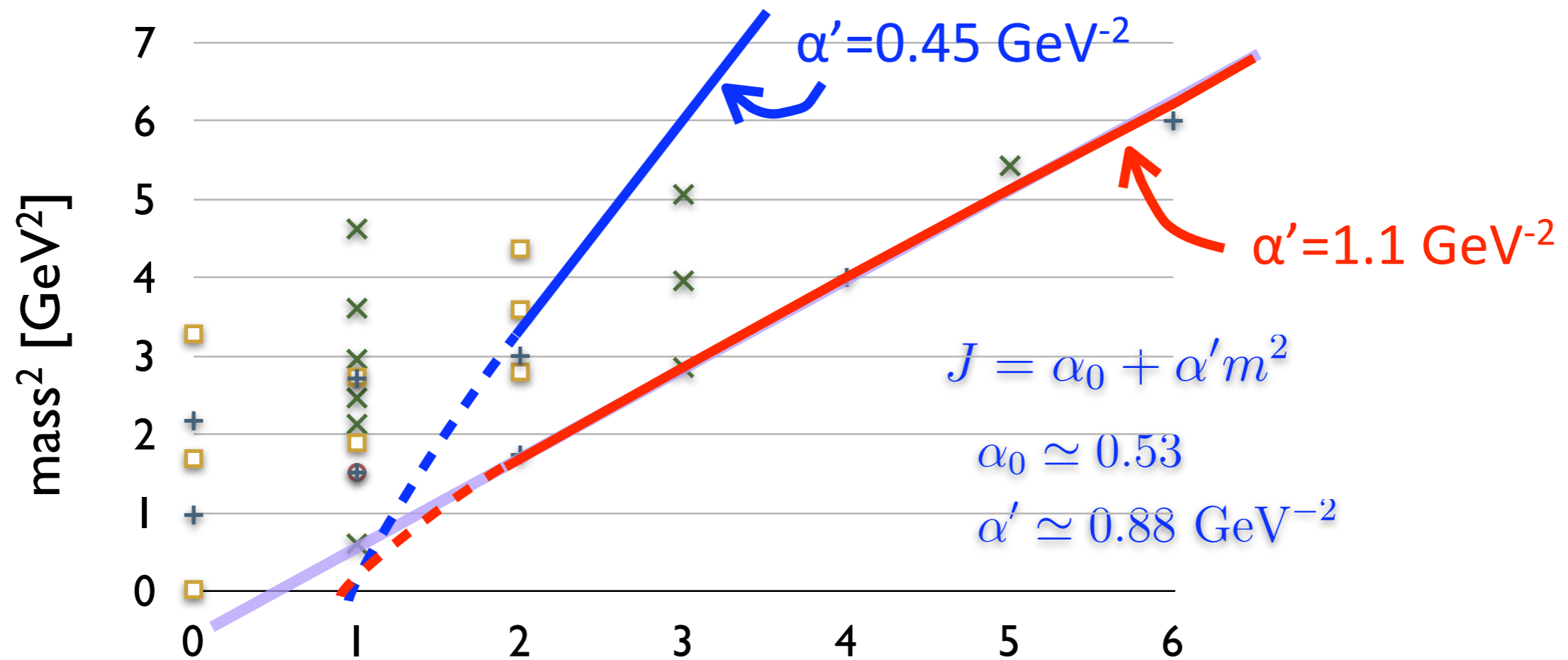
	$B_\mu^{(1)}$	$B_\mu^{(2)}$	$B_\mu^{(3)}$	$B_\mu^{(3)}$...	$\varphi^{(0)}$
J^{PC}	1^{--}	1^{++}	1^{--}	1^{++}	...	0^{-+}
	ρ	a_1	ρ'	a'_1	...	π
mass (MeV)	[776]	1189	1607	2024	...	0

experiment:

$1^{--}(\rho)$	776	1459	1570 Δ	1720	1900 Δ	2150 Δ
$1^{++}(a_1)$	1230	1647 Δ				



Regge trajectory



$$M_n \simeq \sqrt{\frac{N}{\alpha'}} + \frac{1}{\sqrt{2}} \left(n + \frac{1}{2} \right) M_{\text{KK}} \quad \xrightarrow{\text{Spin}} \quad J \simeq 1 + \alpha' M^2 - \frac{\alpha'}{\sqrt{2}} M_{\text{KK}} M + \mathcal{O}(\lambda^{-1})$$

$N = J - 1, n = 0$

- If we use f_π to fit λ , we obtain $\alpha' = 0.45 \text{ GeV}^{-2}$. This is unfortunately too small.
- If we set $\alpha' = 1.1 \text{ GeV}^{-2}$ we get very good fit.

● First excited states ($N=1, n=0$)

J^{PC}	2^{++}	1^{+-}	1^{--}	0^{-+}	$0^{++} \times 3$
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$2^{++}, 1^{+-}, 0^{-+}, 0^{++}$ cannot be $N=0$
 \Rightarrow good candidates for $N=1$

$0^{-+}(\pi)$	135	1300	1812			
$0^{++}(a_0)$	*985	1474				
$1^{--}(\rho)$	776	1459	1570 Δ	1720	1900 Δ	2150 Δ
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$4^{++}(a_4)$	2001					
$5^{--}(\rho_5)$	2330 Δ					
$6^{++}(a_6)$	2450 Δ					

■ : $N=0$
 ■ : $N=1$

- degenerate around 1300 MeV
- * $a_0(980)$ is considered to be a four quark state.

● Second excited states ($N=2, n=0$)

J^{PC}	3^{--}	2^{++}	2^{--}	$2^{-+} \times 2$	$1^{--} \times 7$	$1^{++} \times 3$	$1^{+-} \times 4$	1^{-+}	$0^{++} \times 2$	$0^{-+} \times 6$
$0^{-+} (\pi)$										
$0^{++} (a_0)$										
$1^{--} (\rho)$										
$1^{++} (a_1)$										
$1^{+-} (b_1)$										
$1^{-+} (\pi_1)$										
$2^{++} (a_2)$										
$2^{-+} (\pi_2)$										
$3^{--} (\rho_3)$										
$4^{++} (a_4)$										
$5^{--} (\rho_5)$										
$6^{++} (a_6)$										




$0^{-+} (\pi)$	135	1300	1812			
$0^{++} (a_0)$	985	1474				
$1^{--} (\rho)$	776	1459	1570 [△]	1720	1900 [△]	2150 [△]
$1^{++} (a_1)$	1230	1647 [△]				
$1^{+-} (b_1)$	1230					
$1^{-+} (\pi_1)$	1376 *	1653				
$2^{++} (a_2)$	1318	1732 [△]				
$2^{-+} (\pi_2)$	1672	1895	2090 [△]			
$3^{--} (\rho_3)$	1689	1990 [△]	2250 [△]			
$4^{++} (a_4)$	2001					
$5^{--} (\rho_5)$	2330 [△]					
$6^{++} (a_6)$	2450 [△]					

	: $N=0$
	: $N=1$
	: $N=2$

- ★: prediction ?
- degenerate around 1700 MeV
- * $\pi_1 (1400)$ is claimed to be a four quark state. (could be hybrid)

Summary

$0^{-+}(\pi)$	135	1300	1812			
$0^{++}(a_0)$	985	1474				
$1^{--}(\rho)$	776	1459	1570 Δ	1720	1900 Δ	2150 Δ
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 : $N=0$
 : $N=1$
 : $N=2$

I think this is non-trivial.
What do you think?

5 Discussion

- Mesons are Strings
- Wikipedia says:

Problems and controversy

Although string theory comes from physics, some say that string theory's current untestable status means that it should be classified as more of a mathematical framework for building models as opposed to a physical theory.

..... Yet, for all this activity, not a single new testable prediction has been made, not a single theoretical puzzle has been solved.

**Don't criticize string theory
in this way anymore !**